

Linear momentum

1 particle: define $\underline{p} = m\underline{v}$ so $\frac{d\underline{p}}{dt} = \underline{F}$

just a definition, rescaling of velocity

But: useful for multi-particle situations where different parts of the system move differently

connection to physical symmetries

connection to QM - particles have wave functions

$\psi(\underline{r}, t)$ which distort with no unique speed

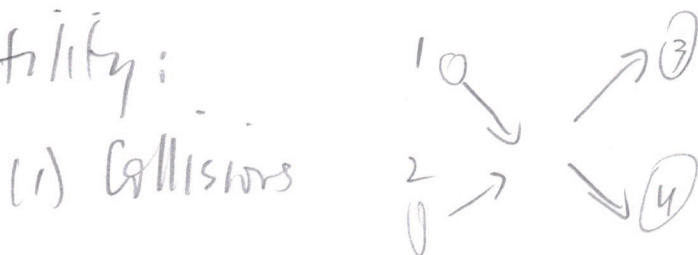
but $p_x \rightarrow -i\hbar \frac{\partial}{\partial x}$ etc. is well-defined

Hamiltonian formalism uses \underline{r} and $\underline{p} \Rightarrow$ QM

Many particles: define $\underline{P} = \sum_{i=1}^N m_i \underline{v}_i = M \underline{v}_{cm}$

$$\text{so } \frac{d\underline{P}}{dt} = \sum_i m_i \underline{\dot{v}}_i = M \frac{d^2 \underline{R}}{dt^2} = \underline{F}^{ext}$$

Utility:



individual \underline{v}_i change but $\underline{F}^{ext} = 0 \Rightarrow \underline{P} = \text{constant}$
useful for analyzing final states

e.g. if $(1) + (2) \rightarrow (3)$ then $\underline{p}_{\text{before}} = \underline{p}_{\text{after}}$

$$\underline{p}_1 + \underline{p}_2 = \underline{p}_3 = m_3 \underline{v}_3$$

so v of product known

likewise if $(A) \rightarrow (B) + (C)$

then in rest frame of A $\underline{p}_A = 0 = \underline{p}_B + \underline{p}_C$

$$\text{so } m_A \underline{v}_A = - m_C \underline{v}_C \quad \text{or } \underline{v}_B = - \underline{v}_C$$

$$\underline{v}_B / \underline{v}_C = m_C / m_B$$

for ≥ 4 particles not that useful,

need to consider cons. of energy also - next chapter

(2) when $\underline{F}^{\text{ext}} = 0$ system is not directed to move in any particular direction

\rightarrow all parts in space are equivalent

\rightarrow "translation invariance" $\hookrightarrow \underline{p} = \text{const}$

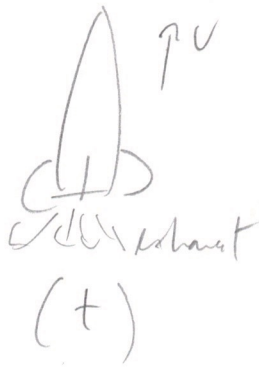
Example of a general relation b/w symmetry & existence of a conserved quantity: more later

(3) If mass added/removed, \underline{p} is what responds to \underline{F}

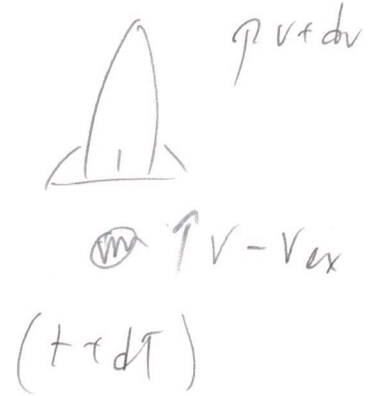
\rightarrow eq. of motion

Rockets

before



after



$$p(t) = m(t) v(t)$$

$$p(t+dt) = \underbrace{(m+dm)}_{dm < 0} \underbrace{(v+dv)}_{dv > 0} + \underbrace{(-dm)}_{\text{pos.}} (v-v_{ex})$$

no gravity

$$= mv + m dv + \cancel{v dm} + \cancel{dm dv} - \cancel{v dm} + v_{ex} dm$$

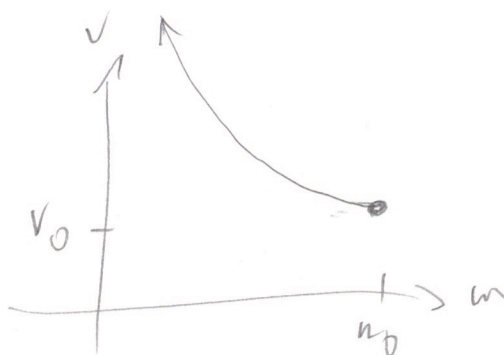
" p(t)

so $m dv = -v_{ex} dm$

$$\frac{v}{v_{ex}} = -\log m + \text{const}$$

$\int \frac{v_0}{v_{ex}} + \log m_0$

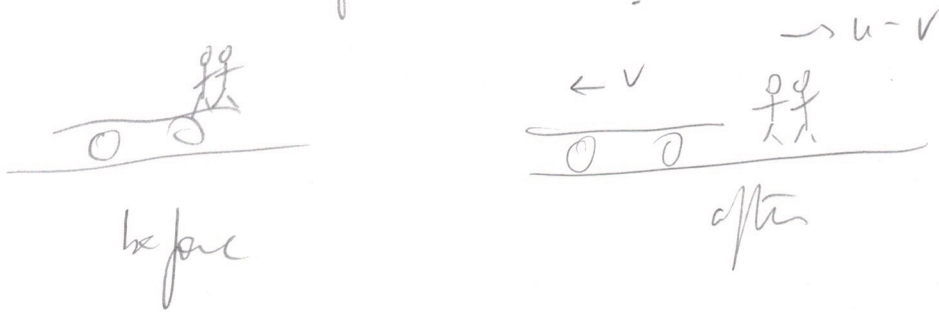
or $v = v_0 + v_{ex} \log \frac{m_0}{m}$



Gravity?
maybe in PS 2

Example (3.4) Two labors (m_h) on an MK car, jump off with relative velocity u . No prob

(a) Both jump off one end simultaneously, car recoils with speed $v = ?$



$$p = 0 = -m_c v + 2m_h(u-v) \quad v = \frac{2m_h u}{m_c + 2m_h}$$

(b) One jumps, then the 2nd:

1st jump: $0 = -(m_c + m_h)v_1 + m_h(u - v_1)$

$$\rightarrow v_1 = \frac{m_h}{m_c + m_h} u$$

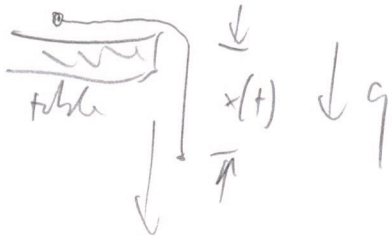
2nd:

$$-(m_c + m_h)v_1 = -m_c v_2 + m_h(u - v_2)$$

$$v_2 = \frac{3m_h + 2m_c}{2m_h + 2m_c} \cdot \frac{2m_h}{2m_h + 2m_c} u < \frac{3m_h + 2m_c}{2m_h + 2m_c} v^{(a)} > v_1 \text{ or } v^{(a)}$$

The labor who jumped first is irrelevant here.

Example: uniform chain of mass m + length l sits on a table, with a segment x_0 hanging over the edge. Gravity will pull the chain off (no friction) - when?



weight of hanging portion $m \frac{x(t)}{l} g$
but chain moves as a unit

$$\text{So } \frac{dy}{dt} = m \ddot{x} = m \frac{x}{l} g \quad \text{or } \ddot{x} = \alpha^2 x \quad \alpha = \sqrt{\frac{g}{l}}$$

$$\rightarrow x(t) = A e^{\alpha t} + B e^{-\alpha t}$$

$$x(0) = x_0 \rightarrow A + B = x_0$$

$$\dot{x}(0) = 0 \rightarrow \alpha(A - B) = 0 \rightarrow A = B = \frac{x_0}{2}$$

$$x(t) = x_0 \cosh \alpha t$$

Chain comes off when $x(t) = l$ or $t = \frac{1}{\alpha} \cosh^{-1} \frac{l}{x_0}$
at which point $v = \alpha x_0 \sinh \alpha t = \alpha x_0 \sinh \left(\cosh^{-1} \frac{l}{x_0} \right)$

To look at this suppose $l \gg x_0$ + use $\cosh z \approx \frac{1}{2} e^z$ + $\sinh z \approx \frac{1}{2} e^z$

$$\Rightarrow v \rightarrow \alpha x_0 \frac{l}{x_0} = \sqrt{gl}$$

If the chain had fallen freely through a distance l , $v = \sqrt{2gl}$
which is larger: less v because less effective g

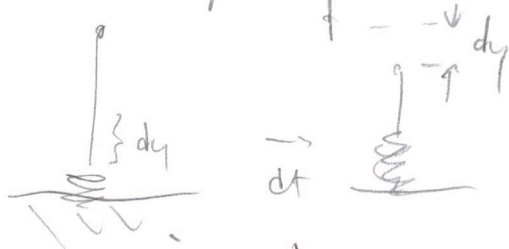
A uniform flexible frictionless chain falls vertically on a table;
 What force does it exert:



$y(t)$ = distance top falls
 = length on the table

$F \neq$ weight on table because falling part transfers momentum.

look at a segment of length dy before & after impact:



momentum transferred = $-\lambda dy \cdot v = dp$

$$F = \text{weight} + \frac{dp}{dt} = -\lambda g y + \frac{(-\lambda dy \cdot v)}{dt} = -\lambda g y - \lambda v^2$$

but for freefall $v^2 = 2gy$

$\rightarrow F = -3\lambda g y$

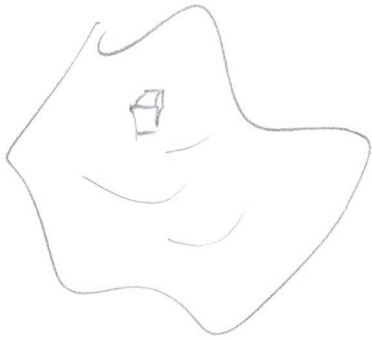
Once the whole chain is on the table $F \rightarrow -\lambda g h$!

difference: moving chain is transferring momentum to the table, more & more as the chain speeds up, but as rest then just gravity.

Continuous systems -

break the continuum into regions ΔV centered at \underline{r}_i

with $m_i = \rho(\underline{r}_i) \Delta V$ $\rho = \text{density}$



$$\underline{R} = \frac{1}{\sum m_i} \sum m_i \underline{r}_i$$

$$\rightarrow \frac{1}{\sum_i \rho(\underline{r}_i) \Delta V} \sum_i \rho(\underline{r}_i) \Delta V \underline{r}_i$$

$$\xrightarrow{\Delta V \rightarrow 0} \frac{1}{\int \rho(\underline{r}) dV} \int \rho(\underline{r}) \underline{r} dV$$

$$= \frac{1}{M} \int \rho(\underline{r}) \underline{r} dV$$

then $\underline{P} = \frac{1}{M} \int \rho(\underline{r}) \underline{r} dV$ etc.

To go back to the continuum:

a point particle at \underline{r}_i has density $= m_i \delta(\underline{r} - \underline{r}_i)$

so $\rho(\underline{r}) = \sum_i m_i \delta(\underline{r} - \underline{r}_i)$ for a set of point particles.

Here " $\delta(\underline{r})$ " = $\delta^3(\underline{r}) = \delta(x) \delta(y) \delta(z)$

$$\delta(x) = \lim_{\epsilon \rightarrow 0} \frac{1}{\sqrt{2\pi\epsilon^2}} e^{-x^2/2\epsilon^2}$$

