

# First Exam Solutions

$$3. \quad L = \sum_{i=1}^N \frac{1}{2} m \dot{\underline{r}}_i^2 - \sum_{i < j} V_2(\underline{r}_i - \underline{r}_j)$$

$$\text{constraint: } \frac{1}{N} \sum \dot{\underline{r}}_i = \underline{v}(t) \rightarrow \sum d\underline{r}_i - N \underline{v} dt = 0$$

method: if  $L(q, \dot{q}, t)$  has  $k$  constraints of the form

$$\sum_i a_{ji}(t) dq_i = 0, \quad j=1, 2, \dots, k \quad \text{the equations of motion are}$$

$$\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{q}_i} - \frac{\partial \mathcal{L}}{\partial q_i} = \sum_{j=1}^k \lambda_j a_{ji}(t)$$

So here

$$\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{\underline{r}}_i} - \frac{\partial \mathcal{L}}{\partial \underline{r}_i} = m \ddot{\underline{r}}_i - \sum_{j \neq i} \underline{F}_{ij} = \underline{\lambda}$$

$$\text{where } \underline{F}_{ij} = - \frac{\partial}{\partial \underline{r}_i} V_2(\underline{r}_i - \underline{r}_j)$$

To find  $\underline{\lambda}$ , add the eqs:

$$\sum m \ddot{\underline{r}}_i + 0 = N \underline{\lambda} \quad (\text{using } \underline{F}_{ij} = -\underline{F}_{ji})$$

$$\text{or } \underline{\lambda} = m \dot{\underline{v}}(t)$$

$$\Rightarrow m \ddot{\underline{r}}_i = \sum \underline{F}_{ij} + m \dot{\underline{v}}(t)$$

$$\text{Check: add the eqs. again: } \sum m \ddot{\underline{r}}_i = m N \dot{\underline{v}}$$

$$\text{or } \frac{d}{dt} \left( \frac{1}{N} \sum \dot{\underline{r}}_i - \underline{v} \right) = 0 \quad ; \quad \text{if constraint obeyed at } t=0 \text{ then it holds for all } t.$$

1. In spherical coordinates, the element of arc length is

$$ds^2 = dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2$$

So on the surface of a sphere ( $r=1$  say), if  $\theta$  &  $\phi$  are written as functions of a parameter  $\lambda$  then

$$D_{12} = \int_1^2 ds = \int_1^2 d\lambda \sqrt{\left(\frac{d\theta}{d\lambda}\right)^2 + \sin^2 \theta \left(\frac{d\phi}{d\lambda}\right)^2}$$

The Euler-Lagrange eqs  $\rightarrow$

$$\frac{d}{d\lambda} \left( \frac{\theta'(\lambda)}{\sqrt{\theta'^2 + \sin^2 \theta \phi'^2}} \right) = \frac{\sin \theta \cos \theta \phi'^2}{\sqrt{\dots}}$$

$$\frac{d}{d\lambda} \left( \frac{\sin^2 \theta \phi'(\lambda)}{\sqrt{\dots}} \right) = 0$$

Without loss of generality the 2 points can be chosen as  $(\theta=0, \phi=0)$  &  $(\theta=\theta_0, \phi=0)$  &

the solution of the 2 eqns subject to these initial conditions is

$$\theta(\lambda) = \lambda \theta_0, \quad 0 \leq \lambda \leq 1, \quad \phi(\lambda) = 0$$

which is part of a great circle.



Method 2: write the curve as  $\underline{r} = \underline{r}(t)$   $t = \text{label of points}$

$$\text{then } ds = \sqrt{d\underline{r}^2(t)} = \sqrt{\underline{r}'(t)^2} dt \quad \& \quad \text{the}$$

problem is to minimize  $\int ds$  subject to  $\underline{r}^2 = 1$  say.

$$\rightarrow \delta \int dt \left[ \sqrt{\underline{r}'(t)^2} + \lambda (\underline{r}^2(t) - 1) \right] = 0$$

$$\text{so } \frac{d}{dt} \left( \frac{\underline{r}'(t)}{\sqrt{\underline{r}'^2(t)}} \right) = \frac{d}{dt} \left( \lambda \underline{r}'(t) \right) = 2\lambda \underline{r}(t)$$

Take  $\hat{r}(t) \times$  this equation:

$$\frac{d}{dt} (\hat{r} \times \hat{r}') = 0 \quad \text{so } \hat{r} \times \hat{r}' = \underline{K} = \text{const. vector.}$$

$$\text{Also } \underline{r}^2 = 1 \rightarrow \underline{r} \cdot \underline{r}' = 0 = \hat{r} \cdot \hat{r}'$$

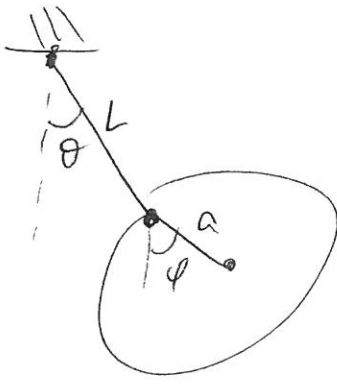
So  $\hat{r}'(t)$  is tangent to the sphere, and  $\perp$  to a fixed space vector  $\underline{K}$ , which is only possible if  $\hat{r}'$  moves on a great circle

e.g. Start at the north pole & assume  $\hat{K} = \text{east}$ .

Then  $\hat{r}'$  must point south at  $t=0$ , and

if it continues south it maintains  $\hat{r} \times \hat{r}' = \underline{K}$ .

2.



$$V = V_{\text{rod}} + V_{\text{disc}}$$

$$= -\frac{mgL}{2} \cos \theta - mg(L \cos \theta + a \cos \phi)$$

$$T = T_{\text{rod}} + T_{\text{disc,cm}} + T_{\text{disc,rel}}$$

$$T_{\text{rod}} = \frac{1}{2} \bar{I}_r \dot{\theta}^2, \quad \bar{I}_r = \text{moment of inertia about an end of the rod}$$

$$T_{\text{disc,cm}} = \frac{m}{2} (\dot{x}_c^2 + \dot{y}_c^2) \quad \text{where } \begin{cases} x_c = L \sin \theta + a \sin \phi \\ y_c = L \cos \theta + a \cos \phi \end{cases}$$

$$= \frac{m}{2} (L^2 \dot{\theta}^2 + a^2 \dot{\phi}^2 + 2aL \dot{\theta} \dot{\phi} \cos(\theta - \phi))$$

$$T_{\text{disc,rel}} = \frac{1}{2} \bar{I}_d \dot{\phi}^2, \quad \bar{I}_d = \text{moment of inertia of disc about center}$$

$$\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{\theta}} - \frac{\partial \mathcal{L}}{\partial \theta} = 0 = \frac{d}{dt} \left( \bar{I}_r \dot{\theta} + mL^2 \dot{\theta} + mL \dot{\phi} \cos(\theta - \phi) \right)$$

$$+ mL \dot{\theta} \dot{\phi} \sin(\theta - \phi) + \frac{mgL}{2} \sin \theta + mgL \sin \theta$$

$$\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{\phi}} - \frac{\partial \mathcal{L}}{\partial \phi} = 0 = \frac{d}{dt} \left( ma^2 \dot{\phi} + \bar{I}_d \dot{\phi} + mL \dot{\theta} \cos(\theta - \phi) \right)$$

$$+ mL \dot{\theta} \dot{\phi} \sin(\theta - \phi) + mga \sin \phi$$

$$0 = (\bar{I}_r + mL^2) \ddot{\theta} + mL \dot{\phi} \cos(\theta - \phi) - mL \dot{\phi}^2 \sin(\theta - \phi) - \frac{3}{2} mgL \sin \theta$$

$$0 = (ma^2 + \bar{I}_d) \ddot{\phi} + mL \ddot{\theta} \cos(\theta - \phi) - mL^2 \dot{\theta}^2 \sin(\theta - \phi) + mga \sin \phi$$

$$4. \quad L = \frac{m}{2} (\dot{r}^2 + r^2 \dot{\theta}^2) - V(r) \quad \rightarrow$$

$$m\ddot{r} = -m r \dot{\theta}^2 - V'(r) \quad \text{and} \quad l = m r^2 \dot{\theta} = \text{constant}$$

$$\text{so} \quad m\ddot{r} = -\frac{l^2}{m r^3} - V'(r)$$

$$\text{Also} \quad \dot{E} = \frac{m}{2} \dot{r}^2 + \frac{l^2}{2mr^2} + V(r) = \frac{m}{2} \dot{r}^2 + V_{\text{eff}}(r) = \text{constant}$$

A <sup>stable</sup> circular orbit at  $r=R$  requires a minimum  $V_{\text{eff}}$ :

$$\frac{\partial V_{\text{eff}}}{\partial r} = 0 \quad \text{and} \quad \frac{\partial^2 V_{\text{eff}}}{\partial r^2} > 0 \quad \text{at} \quad r=R$$

$$\rightarrow V'(R) = \frac{l^2}{mR^3}, \quad V''(R) + \frac{3l}{mR^4} = V''(R) + \frac{3}{R} V'(R) > 0$$

(Also  $\dot{r} = 0$  at  $r=R \rightarrow E = V_{\text{eff}}(R)$  which fixes  $E$ .)

For small oscillations write  $r = R + x$ ,  $|x| \ll R$ , &

$$\begin{aligned} m\ddot{x} &= -\frac{l^2}{m(R+x)^3} - V'(R+x) \\ &= -\frac{l^2}{mR^3} + \frac{3l^2 x}{mR^4} - V'(R) - x V''(R) + O(x^2) \\ &= -\left(V''(R) + \frac{3}{R} V'(R)\right) x = -m\omega^2 x \end{aligned}$$

Note  $\omega^2 > 0$  as required. If  $V(r) = -kr^n$  for example,  
 $V''(R) + \frac{3}{R} V'(R) = -kn R^{n-2} (n+2)$  so  $n > -2$ .